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# Boundary value problems associated with generalized $Q$ -holomorphic functions

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## Abstract

In this work, we discuss Riemann-Hilbert and its adjoint homogeneous problem associated with generalized  $Q$ -holomorphic functions and investigate the solvability of the Riemann-Hilbert problem.

**Keywords:** generalized Beltrami systems;  $Q$ -holomorphic functions; Riemann-Hilbert problem

## Introduction

Douglis [1] and Bojarskii [2] developed an analog of analytic functions for elliptic systems in the plane of the form

$$w_{\bar{z}} - qw_z = 0, \quad (1)$$

where  $w$  is an  $m \times 1$  vector and  $q$  is an  $m \times m$  quasi-diagonal matrix. Also, Bojarskii assumed that all eigenvalues of  $q$  are less than 1. Such systems are natural ones to consider because they arise from the reduction of general elliptic systems in the plane to a standard canonical form. Subsequently Douglis and Bojarskii's theory has been used to study elliptic systems in the form

$$w_{\bar{z}} - qw_z = aw + b\bar{w} + F$$

and the solutions of such equations were called generalized (or pseudo) hyperanalytic functions. Work in this direction appears in [3–5]. These results extend the generalized (or 'pseudo') analytic function theory of Vekua [6] and Bers [7]. Also, classical boundary value problems for analytic functions were extended to generalized hyperanalytic functions. A good survey of the methods encountered in a hyperanalytic case may be found in [8, 9], also see [10].

In [11], Hile noticed that what appears to be the essential property of elliptic systems in the plane for which one can obtain a useful extension of analytic function theory is the self-commuting property of the variable matrix  $Q$ , which means

$$Q(z_1)Q(z_2) = Q(z_2)Q(z_1)$$

for any two points  $z_1, z_2$  in the domain  $G_0$  of  $Q$ . Further, such a  $Q$  matrix cannot be brought into a quasi-diagonal form of Bojarskii by a similarity transformation. So, Hile [11] at-

tempted to extend the results of Douglis and Bojarskii to a wider class of systems in the same form with equation (1). If  $Q(z)$  is self-commuting in  $G_0$  and if  $Q(z)$  has no eigenvalues of magnitude 1 for each  $z$  in  $G_0$ , then Hile called the system (1) the generalized Beltrami system and the solutions of such a system were called *Q-holomorphic functions*. Later in [12, 13], using Vekua and Bers techniques, a function theory is given for the equation

$$Dw + Aw + B\bar{w} = 0, \quad \text{where } D := \frac{\partial}{\partial \bar{z}} - Q(z) \frac{\partial}{\partial z}, \tag{2}$$

where the unknown  $w(z) = \{w_{ij}(z)\}$  is an  $m \times s$  complex matrix,  $Q(z) = \{q_{ij}(z)\}$  is a self-commuting complex matrix with dimension  $m \times m$  and  $q_{k,k-1} \neq 0$  for  $k = 2, \dots, m$ .  $A = \{a_{ij}(z)\}$  and  $B = \{b_{ij}(z)\}$  are commuting with  $Q$ . Solutions of such an equation were called *generalized Q-holomorphic functions*.

In this work, as in a complex case, following Vekua (see [6, pp.228-236]), we investigate the necessary and sufficient condition of solvability of the Riemann-Hilbert problem for equation (2).

### Solvability of Riemann-Hilbert problems

In a regular domain  $G$ , we consider the problem

$$(A) : \begin{cases} L[w] := Dw + Aw + B\bar{w} = F & \text{in } G, \\ \text{Re}(\bar{\lambda}w) = \gamma & \text{on } \partial G. \end{cases} \tag{3}$$

We refer to this problem as boundary value problem (A). Where the unknown  $w(z) = \{w_{ij}(z)\}$  is an  $m \times s$  complex matrix-valued function,  $Q = \{q_{ij}(z)\}$  is a Hölder-continuous function which is a self-commuting matrix with  $m \times m$  and  $q_{k,k-1} \neq 0$  for  $k = 2, \dots, m$ .  $A = \{a_{ij}(z)\}$  and  $B = \{b_{ij}(z)\}$  are commuting with  $Q$ , which is

$$Q(z_1)A(z_2) = A(z_1)Q(z_2), \quad Q(z_1)B(z_2) = B(z_1)Q(z_2).$$

It is assumed, moreover, that  $Q$  is commuting with  $\bar{Q}$  and  $\lambda(z) \in C^1(\Gamma)$  is commuting with  $Q$ , where  $\Gamma = \partial G$ ,  $\lambda\bar{\lambda} = I$ . In respect of the data of problem (A), we also assume that  $A, B$  and  $F \in L^{p,2}(\mathbb{C})$  and  $\gamma \in C_\alpha(\Gamma)$ . If  $F \equiv 0, \gamma \equiv 0$ , we have homogeneous problem ( $\overset{\circ}{A}$ ).

We refer to the adjoint, homogeneous problem (A) as ( $\overset{\circ}{A}$ ); it is given by

$$(\overset{\circ}{A}) : \begin{cases} L'[w'] := Dw' - Aw' - B^*\bar{w}' = 0 & \text{in } G, \\ \text{Re}(\frac{d\phi}{ds}\lambda w') = 0 & \text{on } \Gamma, \end{cases} \tag{4}$$

where  $\phi$  is a generating solution for the generalized Beltrami system ([11, p.109]),  $B^* = \phi_z^{-1}\bar{\phi}_z\bar{B}$ ,  $\frac{d\phi}{ds} := \frac{\partial\phi}{\partial z}\frac{dz}{ds} + \frac{\partial\phi}{\partial \bar{z}}\frac{d\bar{z}}{ds}$  and  $ds$  is the arc length differential. From the Green identity for  $Q$ -holomorphic functions (see [11, p.113]), we have

$$\text{Re} \left[ \frac{1}{2i} \int_\Gamma d\phi w' w - \iint_G \phi_z (w' L[w] - L'[w'] w) dx dy \right] = 0, \tag{5}$$

where  $w'$  is commuting with  $Q$ . For  $L[w] = F$  and  $L'[w'] = 0$ , this becomes

$$\frac{1}{2i} \int_\Gamma d\phi(z) w'(z) \lambda(z) \gamma(z) - \text{Re} \left( \iint_G \phi_z(z) w'(z) F(z) dx dy \right) = 0. \tag{6}$$

Since  $w'$  satisfies the boundary condition

$$\operatorname{Re}\left(\frac{d\phi}{ds}\lambda w'\right) = 0, \tag{7}$$

we have

$$w' = i\lambda^{-1}\left(\frac{d\phi}{ds}\right)^{-1}\varkappa, \tag{8}$$

where  $\varkappa$  is a real matrix commuting with  $Q$ .

The solutions to problem  $(\mathring{A}')$  may be represented by means of fundamental kernels in terms of a real, matrix density  $\varkappa$  as

$$\begin{aligned} w'(z) &= P^{-1} \int_{\Gamma} (d\phi(\zeta)\Omega^{(1)}(z, \zeta)w'(\zeta) - \overline{d\phi(\zeta)}\Omega^{(2)}(z, \zeta)\overline{w'(\zeta)}) \\ &= iP^{-1} \int_{\Gamma} (\Omega^{(1)}(z, \zeta)\lambda^{-1}(\zeta) + \overline{\Omega^{(2)}(z, \zeta)\lambda^{-1}(\zeta)})\varkappa(\zeta) ds, \end{aligned} \tag{9}$$

see ([14, p.543]). In (9),  $P$  is a constant matrix defined by

$$P(z) = \int_{|z|=1} (zI + \bar{z}Q)^{-1}(I dz + Q d\bar{z})$$

called  $P$ -value for the generalized Beltrami system [11]. Since  $\varkappa$  is a real matrix commuting with  $Q$ , inserting the expression (9) into the boundary condition (7), we have

$$\int_{\Gamma} K_1(\zeta, z)\varkappa(\zeta) ds_{\zeta} = 0, \quad z, \zeta = \zeta(s) \in \Gamma, \tag{10}$$

where

$$K_1(\zeta, z) = -\operatorname{Re}\left[ iP^{-1}\lambda(z)\frac{d\phi(z)}{ds}(\Omega^{(1)}(z, \zeta)\lambda^{-1}(z) + \overline{\Omega^{(2)}(z, \zeta)\lambda^{-1}(z)}) \right].$$

The integral in (10) is to be taken in the Cauchy principal value sense. If we denote this equation in an operator form by  $\tilde{K}\varkappa = 0$  and its adjoint by  $\tilde{K}'f = 0$ , then it may be easily demonstrated that the index of (10) is  $\kappa = k - k' = 0$ . Here  $k$  and  $k'$  are dimensions of null spaces of  $\tilde{K}$  and  $\tilde{K}'$  respectively. If  $\{\varkappa_1, \dots, \varkappa_k\}$  is a complete system of solutions of (10), putting each of this into (9), we obtain the solutions of problem  $(\mathring{A}')$ . However, it is possible that some of these solutions may turn out to be trivial solutions, which occurs when  $(\lambda\frac{d\phi}{ds})^{-1}\varkappa$  takes on the boundary values of a  $Q$ -holomorphic function  $\psi_j$  on each component of boundary contours  $\Gamma_j$  in  $C^{1,\alpha}(\mathbb{C})$  which is, moreover,  $Q$ -holomorphic in the domain  $G_j$  bounded by the closed contour  $\Gamma_j$ . Let  $\{\varkappa_1, \dots, \varkappa_{\ell'}\}$  be solutions of equation (10) to which linearly independent solutions (see [15])  $w'_1, \dots, w'_{\ell'}$  of problem  $(\mathring{A}')$  correspond, then the remaining solutions  $\{\varkappa_{\ell'+1}, \dots, \varkappa_k\}$  satisfy the boundary condition of the form

$$\varkappa(z) = i\lambda(z)\frac{d\phi}{ds}\Phi^-(z) \quad \text{on } \Gamma. \tag{11}$$

Here  $\Phi^-$  are meant to be  $Q$ -holomorphic functions outside of  $G^- := G + \Gamma$  and  $\Phi(\infty) = 0$ . Hence the  $Q$ -holomorphic functions satisfy the homogeneous boundary conditions

$$\mathring{A}'_* := \operatorname{Re}\left(\lambda \frac{d\phi}{ds} \Phi^-\right) = 0 \quad \text{on } \Gamma. \tag{12}$$

In a complex case, Vekua refers to problems of this type as being concomitant to  $(\mathring{A}')$  and denotes them by  $(\mathring{A}'_*)$ . Let  $\ell_*$  be a number of linearly independent solutions of this problem. Obviously,  $\ell' + \ell'_* = k$ .

Let us now return to the discussion of problem  $(A)$ , where we assume that  $\varkappa = 0$  in what follows. The solutions of this problem may be expressed in terms of the generalized Cauchy kernel as follows:

$$w(z) = w_1(z) + w_2(z) = \underset{\sim}{C}[\lambda\gamma](z) + \underset{\sim}{C}[i\lambda\mu](z),$$

where

$$\underset{\sim}{C}[\Phi] = P^{-1} \int_{\Gamma} d\phi(\zeta) \Omega^{(1)}(z, \zeta) \Phi(\zeta) - \overline{d\phi(\zeta)} \Omega^{(2)}(z, \zeta) \overline{\Phi(\zeta)}$$

(see [14, p.543]). From the Plemelj formulas, it is seen that the density  $\mu$  must satisfy the integral equation

$$\gamma_0 = \int_{\Gamma} K_1(\zeta, z) \mu(z) ds_z, \tag{13}$$

where

$$\gamma_0(\zeta) = \gamma(\zeta) - \operatorname{Re}[\overline{\lambda(\zeta)} w_1^+(\zeta)] = -\operatorname{Re}[\overline{\lambda(\zeta)} w_1^-(\zeta)]. \tag{14}$$

Problem  $(\mathring{A}'_*)$  concomitant to problem  $(\mathring{A}')$  has the boundary condition  $\operatorname{Re}[\lambda^{-1}(z) \Phi^-(z)] = 0$  on  $\Gamma$ , where  $\Phi$  is  $Q$ -holomorphic outside  $G + \Gamma$  and  $\Phi(\infty) = 0$ . Denoting the numbers of linearly independent solutions of  $(\mathring{A}')$  and  $(\mathring{A}'_*)$  by  $\ell$  and  $\ell_*$  respectively, we have  $k = \ell + \ell_*$ . In order that (13) is solvable, it is necessary and sufficient that the nonhomogeneous data  $\gamma_0$  satisfy the auxiliary conditions

$$\int_{\Gamma} \varkappa_j(\zeta) \gamma_0(\zeta) ds_{\zeta} = 0 \quad (j = 1, \dots, k), \tag{15}$$

where  $\varkappa_j$  are solutions to integral equation (10). These solutions may be broken up into two groups  $\{\varkappa_1, \dots, \varkappa_{\ell'}\}$  and  $\{\varkappa_{\ell'+1}, \dots, \varkappa_k\}$  such that  $\varkappa_j = i\lambda(z) \frac{d\phi}{ds} w'_j(z)$  for  $j = 1, \dots, \ell'$  and  $\varkappa_j = i\lambda(z) \frac{d\phi}{ds} \Phi_j$  for  $j = \ell' + 1, \dots, k$ , where  $z \in \Gamma$ . Here  $w'_j$  and  $\Phi_j$  are solutions of problems  $(\mathring{A}')$  and  $(\mathring{A}'_*)$  respectively. The condition (15) for  $\gamma_0$  given by (14) becomes

$$-\int_{\Gamma} \varkappa_j(\zeta) \gamma_0(\zeta) ds_{\zeta} = -i \int_{\Gamma} d\phi(\zeta) \lambda(\zeta) w'_j(\zeta) \gamma(\zeta) + \operatorname{Re}\left[ i \int_{\Gamma} d\phi(\zeta) w'_j(\zeta) w_1^+(\zeta) \right]$$

for  $j = 1, \dots, \ell'$ , whereas for  $j = \ell' + 1, \dots, k$ , we have

$$\int_{\Gamma} \varkappa_j(\zeta) \gamma_0(\zeta) ds = \operatorname{Re}\left[ i \int_{\Gamma} d\phi(\zeta) \Phi_j^-(\zeta) w_1^-(\zeta) \right] = 0.$$

Consequently, the conditions (15) are seen to hold if (6) (with  $F = 0$ ) holds. From the above discussion, one obtains a Fredholm-type theorem for problem (A).

**Theorem 1** *Non-homogeneous boundary problem (A) is solvable if and only if the condition (6) is satisfied,  $w'$  being an arbitrary solution of adjoint homogeneous boundary problem ( $\dot{A}'$ ).*

This theorem immediately implies the following.

**Theorem 2** *Non-homogeneous boundary problem (A) is solvable for an arbitrary right-hand side if and only if adjoint homogeneous problem ( $\dot{A}'$ ) has no solution.*

#### Competing interests

The author declares that they have no competing interests.

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